

Operator Methods Applied to Special Functions

Francisco Soto-Eguibar and Héctor M. Moya-Cessa*

Instituto Nacional de Astrofísica, Óptica y Electrónica, INAOE. Luis Enrique Erro 1, Santa María Tonantzintla, San Andrés Cholula, Puebla, 72840, México

Received: 23 Jul. 2014, Revised: 24 Oct. 2014, Accepted: 25 Oct. 2014
Published online: 1 May 2015

Abstract: Based on operator algebras commonly used in quantum mechanics some properties of special functions such as Hermite and Laguerre polynomials and Bessel functions are derived.

Keywords: operator algebras, special functions, Hermite polynomials, Laguerre polynomials, Bessel functions

1 Introduction

Based on some techniques that are common to quantum mechanics, we give some examples on how some series of special functions may be added.

For this, we use some theorems and lemmas that are not usually known when one studies special functions, such as the Baker-Hausdorff formula and the Hadamard lemma. We developed further preliminary results presented in [1].

2 Hermite polynomials

The generating function for the Hermite polynomials [2] is [3,4,5]

$$e^{-\alpha^2+2\alpha x} = \sum_{n=0}^{\infty} H_n(x) \frac{\alpha^n}{n!}. \quad (1)$$

The Hermite polynomials may be obtained from Rodrigues' formula [3,4,6,7] as

$$H_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} e^{-x^2}. \quad (2)$$

From the recurrence relations [3]

$$H_{n+1}(x) = 2xH_n(x) - 2nH_{n-1}(x), \quad (3)$$

and from

$$\frac{dH_n(x)}{dx} = 2nH_{n-1}(x), \quad (4)$$

we can generate all the Hermite polynomials. From the above recurrence relations, we can also prove

that, if we define the functions

$$\psi_n(x) = \frac{\pi^{-1/4}}{\sqrt{2^n n!}} e^{-x^2/2} H_n(x), \quad n = 0, 1, 2, \dots, \quad (5)$$

then

$$\hat{A}^\dagger \psi_n(x) \equiv \frac{1}{\sqrt{2}} \left(x - \frac{d}{dx} \right) \psi_n(x) = \sqrt{n+1} \psi_{n+1}(x) \quad (6)$$

and

$$\hat{A} \psi_n(x) \equiv \frac{1}{\sqrt{2}} \left(x + \frac{d}{dx} \right) \psi_n(x) = \sqrt{n} \psi_{n-1}(x). \quad (7)$$

The functions (5) constitute a complete orthonormal set for the space of square integrable functions, then we can expand any function in that space as

$$f(x) = \sum_{n=0}^{\infty} c_n \psi_n(x), \quad (8)$$

where

$$c_n = \int_{-\infty}^{\infty} dx f(x) \psi_n(x). \quad (9)$$

Hermite polynomials are also solutions of the second order ordinary differential equation [3,4,5]

$$y'' - 2xy' + 2ny = 0. \quad (10)$$

Let us define the differential operator \hat{p} as

$$\hat{p} = -i \frac{d}{dx}, \quad (11)$$

* Corresponding author e-mail: hmmc@inaoep.mx

then we have that

$$\frac{d^n}{dx^n} = \left(-\frac{\hat{p}}{i}\right)^n, \quad (12)$$

and we can rewrite (2) in the form

$$H_n(x) = (-i)^n e^{x^2} \hat{p}^n e^{-x^2} = (-i)^n \left(e^{x^2} \hat{p} e^{-x^2}\right)^n. \quad (13)$$

The operator inside the parenthesis above has the form

$$e^{\xi \hat{A}} \hat{B} e^{-\xi \hat{A}}, \quad (14)$$

for which we will use the following lemma [8, 9, 10]:

Hadamard lemma. Given two linear operators \hat{A} and \hat{B} then

$$\begin{aligned} e^{\xi \hat{A}} \hat{B} e^{-\xi \hat{A}} &= \hat{B} + \xi [\hat{A}, \hat{B}] + \frac{\xi^2}{2!} [\hat{A}, [\hat{A}, \hat{B}]] \\ &+ \frac{\xi^3}{3!} [\hat{A}, [\hat{A}, [\hat{A}, \hat{B}]]] + \dots, \end{aligned} \quad (15)$$

where $[\hat{A}, \hat{B}] \equiv \hat{A}\hat{B} - \hat{B}\hat{A}$ is the commutator of operators \hat{A} and \hat{B} .

This allows us to obtain an expression for the formula (13) developed above. We identify $\xi = 1$, $\hat{A} = x^2$ and $\hat{B} = \hat{p}$ in equation (15), so that

$$\begin{aligned} e^{x^2} \hat{p} e^{-x^2} &= \hat{p} + 1 [x^2, \hat{p}] + \frac{1}{2!} [x^2, [x^2, \hat{p}]] \\ &+ \frac{1}{3!} [x^2, [x^2, [x^2, \hat{p}]]] + \dots \end{aligned} \quad (16)$$

To calculate the first commutator $[x^2, \hat{p}]$, we use the general property $[\hat{A}\hat{B}, \hat{C}] = \hat{A}[\hat{B}, \hat{C}] + [\hat{A}, \hat{C}]\hat{B}$ of commutators, and that $[x, \hat{p}] = i$, to get the commutation relation $[x^2, \hat{p}] = 2ix$. It is obvious that all the other commutators in (16) are zero, and we finally get

$$H_n(x) = (-i)^n (\hat{p} + 2ix)^n 1. \quad (17)$$

This last expression can be used to obtain the generating function. We have,

$$\sum_{n=0}^{\infty} H_n(x) \frac{\alpha^n}{n!} = \sum_{n=0}^{\infty} \frac{\alpha^n}{n!} (-i)^n (\hat{p} + 2ix)^n 1 = e^{-i\alpha(\hat{p}+2ix)} 1. \quad (18)$$

We obtained the exponential of the sum of two quantities that do not commute. The above exponential can be factorized in the product of exponentials via the Baker-Hausdorff formula:

2.1 Baker-Hausdorff formula

Baker-Hausdorff formula [8, 10, 11]. Given two operators \hat{A} and \hat{B} that obey

$$[[\hat{A}, \hat{B}], \hat{A}] = [[\hat{A}, \hat{B}], \hat{B}] = 0, \quad (19)$$

then

$$e^{\hat{A}+\hat{B}} = e^{-\frac{1}{2}[\hat{A}, \hat{B}]} e^{\hat{A}} e^{\hat{B}}. \quad (20)$$

We apply this formula to expression (18) identifying $\hat{A} = 2\alpha x$, $\hat{B} = -i\alpha \hat{p}$, and we get $[\hat{A}, \hat{B}] = -2i\alpha^2 [x, \hat{p}] = 2\alpha^2$, such that

$$\sum_{n=0}^{\infty} H_n(x) \frac{\alpha^n}{n!} = e^{-\alpha^2} e^{2\alpha x} e^{-i\alpha \hat{p}} 1. \quad (21)$$

Using now the obvious fact that $e^{-i\alpha \hat{p}} 1 = 1$, we finally obtain

$$\sum_{n=0}^{\infty} H_n(x) \frac{\alpha^n}{n!} = e^{-\alpha^2 + 2\alpha x}, \quad (22)$$

that is the generating function for Hermite polynomials [12].

2.2 Series of even Hermite polynomials

In order to show the power of the operator methods, we calculate now the value of the following even Hermite polynomials series,

$$F(t) = \sum_{n=0}^{\infty} \frac{t^n}{n!} H_{2n}(x). \quad (23)$$

From (17), we get

$$H_{2n}(x) = (-1)^n (\hat{p} + 2ix)^{2n} 1. \quad (24)$$

Therefore,

$$\begin{aligned} F(t) &= \sum_{n=0}^{\infty} \frac{t^n}{n!} H_{2n}(x) = \sum_{n=0}^{\infty} \frac{t^n}{n!} (-1)^n (\hat{p} + 2ix)^{2n} 1 \\ &= \sum_{n=0}^{\infty} \frac{t^n}{n!} (-1)^n [(\hat{p} + 2ix)^2]^n 1 \\ &= \exp[-t(\hat{p} + 2ix)^2] 1. \end{aligned} \quad (25)$$

Developing the power in the exponential above we get

$$\begin{aligned} F(t) &= \sum_{n=0}^{\infty} \frac{t^n}{n!} H_{2n}(x) \\ &= \exp\{-t[\hat{p}^2 - 4x^2 + 2i(x\hat{p} + \hat{p}x)]\} 1. \end{aligned} \quad (26)$$

The operators in the exponential in this last expression do not satisfy the conditions of the Baker-Hausdorff formula, so we need another method to understand the action of the full operator that appears in the right side of expression (26). What we do is to propose the ansatz,

$$F(t) = \exp[f(t)x^2] \exp[g(t)(x\hat{p} + \hat{p}x)] \exp[h(t)\hat{p}^2] 1, \quad (27)$$

where $f(t)$, $g(t)$ and $h(t)$ are functions to be determined. The exponential that contains the term $g(t)$ is the

so-called squeeze operator [13]. Deriving this expression with respect to t , and dropping the explicit dependence of $F(t)$, $f(t)$, $g(t)$ and $h(t)$ on t ,

$$\begin{aligned} \frac{dF}{dt} &= \frac{df}{dt}x^2F \\ &+ \frac{dg}{dt} \exp(fx^2) (x\hat{p} + \hat{p}x) \exp[g(x\hat{p} + \hat{p}x)] \exp(hp^2) 1 \\ &+ \frac{dh}{dt} \exp(fx^2) \exp[g(x\hat{p} + \hat{p}x)] \hat{p}^2 \exp(hp^2) 1. \end{aligned} \quad (28)$$

Introducing an "smart" 1 in the second and third term, we get

$$\begin{aligned} \frac{dF}{dt} &= \frac{df}{dt}x^2F + \frac{dg}{dt}e^{fx^2} (x\hat{p} + \hat{p}x) e^{-fx^2} F \\ &+ \frac{dh}{dt}e^{fx^2} \exp[g(x\hat{p} + \hat{p}x)] \hat{p}^2 \exp[-g(x\hat{p} + \hat{p}x)] e^{-fx^2} F. \end{aligned} \quad (29)$$

We work then with the operator in the second term; we use the Hadamard lemma (15) to obtain

$$\begin{aligned} e^{fx^2} (x\hat{p} + \hat{p}x) e^{-fx^2} &= x\hat{p} + \hat{p}x + f[x^2, x\hat{p} + \hat{p}x] \\ &+ \frac{f^2}{2!} [x^2, [x^2, x\hat{p} + \hat{p}x]] \\ &+ \frac{f^3}{3!} [x^2, [x^2, [x^2, x\hat{p} + \hat{p}x]]] + \dots \end{aligned} \quad (30)$$

The first commutator that appears in the above expression is easily calculated, $[x^2, x\hat{p} + \hat{p}x] = 4ix^2$, and so all the others commutators are zero. Substituting back in (30), we get

$$\exp(fx^2) (x\hat{p} + \hat{p}x) \exp(-fx^2) = x\hat{p} + \hat{p}x + 4ifx^2. \quad (31)$$

We analyze now the third operator in expression (29). We study first only a part of it, $\exp[g(x\hat{p} + \hat{p}x)] \hat{p}^2 \exp[-g(x\hat{p} + \hat{p}x)]$. Using again (15),

$$\begin{aligned} \exp[g(x\hat{p} + \hat{p}x)] \hat{p}^2 \exp[-g(x\hat{p} + \hat{p}x)] &= \\ &= \hat{p}^2 + g[x\hat{p} + \hat{p}x, \hat{p}^2] + \frac{g^2}{2!} [x\hat{p} + \hat{p}x, [x\hat{p} + \hat{p}x, \hat{p}^2]] \\ &+ \frac{g^3}{3!} [x\hat{p} + \hat{p}x, [x\hat{p} + \hat{p}x, [x\hat{p} + \hat{p}x, \hat{p}^2]]] + \dots \end{aligned} \quad (32)$$

Calculating the first commutators,

$$[x\hat{p} + \hat{p}x, \hat{p}^2] = 4ip^2, \quad (33)$$

$$[x\hat{p} + \hat{p}x, [x\hat{p} + \hat{p}x, \hat{p}^2]] = -16p^2, \quad (34)$$

$$[x\hat{p} + \hat{p}x, [x\hat{p} + \hat{p}x, [x\hat{p} + \hat{p}x, \hat{p}^2]]] = -64ip^2, \quad (35)$$

it is clear that

$$\begin{aligned} \exp[g(x\hat{p} + \hat{p}x)] \hat{p}^2 \exp[-g(x\hat{p} + \hat{p}x)] &= \\ &= \hat{p}^2 \sum_{j=0}^{\infty} (4i)^j \frac{g^j}{j!} = \hat{p}^2 \exp(4ig). \end{aligned} \quad (36)$$

We proceed now to complete the study of the third operator in expression (29). Until now we have

$$\begin{aligned} e^{-fx^2} \exp[g(x\hat{p} + \hat{p}x)] \hat{p}^2 \exp[-g(x\hat{p} + \hat{p}x)] e^{-fx^2} &= \\ &= \exp(-fx^2) \hat{p}^2 \exp(4ig) \exp(-fx^2). \end{aligned}$$

We use once more formula (15), to write

$$\begin{aligned} \exp(fx^2) \hat{p}^2 \exp(-fx^2) &= \hat{p}^2 + f[x^2, \hat{p}^2] \\ &+ \frac{f^2}{2!} [x^2, [x^2, \hat{p}^2]] + \frac{f^3}{3!} [x^2, [x^2, [x^2, \hat{p}^2]]] + \dots \end{aligned} \quad (37)$$

The first commutator gives $[x^2, \hat{p}^2] = -2 + 4ipx$, the second one gives $[x^2, [x^2, \hat{p}^2]] = -8x^2$, and the third one $[x^2, [x^2, [x^2, \hat{p}^2]]] = 0$; such that all the other commutators are zero, and

$$\begin{aligned} \exp(fx^2) \hat{p}^2 \exp(-fx^2) &= \hat{p}^2 + f(-2 + 4ipx) \\ &+ \frac{f^2}{2!} (-8x^2) = \hat{p}^2 + 2if(x\hat{p} + \hat{p}x) - 4f^2x^2. \end{aligned} \quad (38)$$

Finally, we can write a reduced expression for the derivative of the original series $F(t)$ as

$$\begin{aligned} \frac{dF}{dt} &= \\ &= \left\{ \left[\frac{df}{dt} + 4if \frac{dg}{dt} - 4f^2 \exp(4ig) \frac{dh}{dt} \right] x^2 + \right. \\ &\left. \left[\frac{dg}{dt} + 2if \exp(4ig) \frac{dh}{dt} \right] (x\hat{p} + \hat{p}x) + \exp(4ig) \frac{dh}{dt} \hat{p}^2 \right\} F. \end{aligned} \quad (39)$$

We get back to the original expression for the operator, expression (27), and taken the derivative with respect to t ,

$$\begin{aligned} \frac{dF(t)}{dt} &= -[\hat{p}^2 - 4x^2 + 2i(x\hat{p} + \hat{p}x)] \\ &\times \exp\{-t[\hat{p}^2 - 4x^2 + 2i(x\hat{p} + \hat{p}x)]\} 1 \\ &= [-\hat{p}^2 + 4x^2 - 2i(x\hat{p} + \hat{p}x)] F. \end{aligned} \quad (40)$$

Comparing (39) and (40), we get the system of differential equations

$$\begin{aligned} \frac{df}{dt} + 4if \frac{dg}{dt} - 4f^2 \exp(4ig) \frac{dh}{dt} &= 4, \\ \frac{dg}{dt} + 2if \exp(4ig) \frac{dh}{dt} &= -2i, \\ \exp(4ig) \frac{dh}{dt} &= -1. \end{aligned} \quad (41)$$

The initial conditions that we must set on these equations are $f(0) = g(0) = h(0) = 0$, so for $t = 0$ the operator in the right side of (26) is the identity. The solutions then are simply

$$f = \frac{4t}{4t+1}, \quad (42)$$

$$g = -\frac{i}{2} \ln(4t+1), \quad (43)$$

$$h = -\frac{t}{4t+1}. \quad (44)$$

We calculate now explicitly (27). As $\hat{p} = -i\frac{d}{dx}$, it is very easy to see that $\exp(h\hat{p}^2)1 = 1$ and then

$$F(t) = \exp(fx^2) \exp[g(x\hat{p} + \hat{p}x)]1. \quad (45)$$

Using now that $[x, \hat{p}] = i$, we write

$$\begin{aligned} \exp[g(x\hat{p} + \hat{p}x)]1 &= \exp[g(2x\hat{p} - i)]1 \\ &= \exp(-ig) \exp(2gx\hat{p})1 \end{aligned} \quad (46)$$

and it is also clear that

$$\begin{aligned} \exp(2gx\hat{p})1 &= \left[1 + gx\hat{p} + \frac{g^2}{2}(x\hat{p})^2 + \dots \right]1 \\ &= 1 + gx\hat{p}1 + \frac{g^2}{2}(x\hat{p})^2 1 + \dots = 1 \end{aligned} \quad (47)$$

and also that

$$\exp[g(x\hat{p} + \hat{p}x)]1 = \exp(-ig). \quad (48)$$

We then have

$$F(t) = \exp(fx^2 - ig). \quad (49)$$

Substituting the functions f y g ,

$$F(t) = \frac{1}{\sqrt{4t+1}} \exp\left(\frac{4tx^2}{4t+1}\right), \quad (50)$$

and finally, getting back to (26), we get the formula we were looking for

$$\sum_{n=0}^{\infty} \frac{t^n}{n!} H_{2n}(x) = \frac{1}{\sqrt{4t+1}} \exp\left(\frac{4tx^2}{4t+1}\right). \quad (51)$$

2.3 Addition formula

We want to apply the form obtained in (17), $H_n(x) = (-i)^n (\hat{p} + 2ix)^n 1$, to evaluate the quantity $H_n(x+y)$. We write it as

$$H_n(x+y) = (-i)^n \left[-i \frac{d}{d(x+y)} + 2i(x+y) \right]^n, \quad (52)$$

by using the chain rule we have $\frac{d}{d(x+y)} = \frac{1}{2} \left(\frac{\partial}{\partial x} + \frac{\partial}{\partial y} \right)$, so that we may re-express (52) in the form

$$\begin{aligned} H_n(x+y) &= \\ &= \left(\frac{-i}{\sqrt{2}} \right)^n \left(-i \frac{\partial}{\partial \sqrt{2}x} + 2i\sqrt{2}x - i \frac{\partial}{\partial \sqrt{2}y} + 2i\sqrt{2}y \right)^n. \end{aligned} \quad (53)$$

By defining $\hat{p}_X = -i\frac{\partial}{\partial X}$ and $\hat{p}_Y = -i\frac{\partial}{\partial Y}$ with $X = \sqrt{2}x$ and $Y = \sqrt{2}y$, we obtain

$$\begin{aligned} H_n(x+y) &= \\ &= \frac{1}{2^{n/2}} \sum_{k=0}^n \binom{n}{k} (-i)^k (p_X + 2iX)^k (-i)^{n-k} (p_Y + 2iY)^{n-k}, \end{aligned} \quad (54)$$

that by using $H_j(x) = (-i)^n (\hat{p} + 2ix)^j 1$ adds to

$$H_n(x+y) = \frac{1}{2^{n/2}} \sum_{k=0}^n \binom{n}{k} H_k(\sqrt{2}x) H_{n-k}(\sqrt{2}y), \quad (55)$$

which is the addition formula we were looking for.

3 Associated Laguerre polynomials

The generating function for the associated Laguerre polynomials is [3, 12, 14]

$$\sum_{n=0}^{\infty} L_n^\alpha(x) t^n = \frac{1}{(1-t)^{\alpha+1}} \exp\left(\frac{-xt}{1-t}\right), \quad |t| < 1. \quad (56)$$

The associated Laguerre polynomials may be obtained from the corresponding Rodrigues' formula [3, 4, 12, 14]

$$L_n^\alpha(x) = \frac{1}{n!} x^{-\alpha} e^x \frac{d^n}{dx^n} (e^{-x} x^{n+\alpha}). \quad (57)$$

The associated Laguerre polynomials satisfy several recurrence relations. One very useful, when extracting properties of the wave functions of the hydrogen atom, is

$$(n+1)L_{n+1}^\alpha(x) = (2n+\alpha+1-x)L_n^\alpha(x) - (n+\alpha)L_{n-1}^\alpha(x). \quad (58)$$

We will use the operator method outlined above for the Hermite polynomials, to derive the usual explicit expression for the associated Laguerre polynomials. We rewrite expression (57) as

$$L_n^\alpha(x) = \frac{1}{n!} x^{-\alpha} e^x (i\hat{p})^n e^{-x} x^{n+\alpha}, \quad (59)$$

where again the operator $\hat{p} = -id/dx$, defined in (11), is used.

We notice that $e^x (i\hat{p})^n e^{-x} = [e^x (i\hat{p}) e^{-x}]^n$ and that, using (15), $e^x \hat{p} e^{-x} = (\hat{p} + i)$, so

$$L_n^\alpha(x) = \frac{1}{n!} x^{-\alpha} \left(\frac{d}{dx} - 1 \right)^n x^{n+\alpha}. \quad (60)$$

Using the binomial expansion and because $\frac{d^m}{dx^m} x^{n+\alpha} = \frac{(n+\alpha)!}{(n+\alpha-m)!} x^{n+\alpha-m}$, we obtain the usual form for associated Laguerre polynomials,

$$L_n^\alpha(x) = \sum_{k=0}^n \binom{n+\alpha}{n-k} (-1)^k \frac{x^k}{k!}. \tag{61}$$

4 Bessel functions of the first kind of integer order

Bessel functions of the first kind of integer order, $J_n(x)$, are solutions of the Bessel differential equation [15, 16, 17]

$$x^2 y'' + xy' + (x^2 - n^2)y = 0, \tag{62}$$

where n is an integer. They may be obtained from the generating function [3, 15, 16, 17]

$$\exp\left[\frac{x}{2}\left(t - \frac{1}{t}\right)\right] = \sum_{n=-\infty}^{\infty} t^n J_n(x), \tag{63}$$

and also from the following recurrence relations [3, 15, 16, 17]

$$\frac{2n}{x} J_n(x) = J_{n-1}(x) + J_{n+1}(x). \tag{64}$$

Bessel functions of the first kind of integer order may be written as

$$J_n(x) = \sum_{m=0}^{\infty} \frac{(-1)^m x^{2m+n}}{2^{2m+n} m! (m+n)!}, \tag{65}$$

and also the following integral representation is very useful [3]

$$J_n(x) = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-i(n\tau - x \sin \tau)} d\tau. \tag{66}$$

Some other important relations for the Bessel functions of the first kind are the Jacobi-Anger expansions[3]:

$$e^{ix \cos y} = \sum_{n=-\infty}^{\infty} i^n J_n(x) e^{iny} \tag{67}$$

and

$$e^{ix \sin y} = \sum_{n=-\infty}^{\infty} J_n(x) e^{iny}. \tag{68}$$

4.1 Addition formula

Using the operator methods developed in previous sections, we will obtain here the addition formula for the Bessel functions of the first kind of integer order. First, we will derive the following expression for any "well behaved" function f ,

$$f(x+y) = e^{iy \hat{p}_x} f(x) e^{-iy \hat{p}_x} 1, \tag{69}$$

where $\hat{p}_x = -id/dx$ is the operator introduced in expression (11). Because $e^{-iy \hat{p}_x} 1 = 1$, developing the f function in a Taylor series (we call c_n to the coefficients in the expansion) and using the linearity of the $e^{iy \hat{p}_x}$ operator,

$$\begin{aligned} e^{iy \hat{p}_x} f(x) e^{-iy \hat{p}_x} 1 &= e^{iy \hat{p}_x} f(x) = e^{iy \hat{p}_x} \sum_{k=0}^{\infty} c_k x^k \\ &= \sum_{k=0}^{\infty} c_k e^{iy \hat{p}_x} x^k. \end{aligned} \tag{70}$$

Now

$$\begin{aligned} e^{iy \hat{p}_x} x^k &= \sum_{l=0}^{\infty} \frac{(iy)^l}{l!} (-i)^l \frac{d^l}{dx^l} x^k = \sum_{l=0}^k (y)^l \binom{k}{l} x^{k-l} \\ &= (x+y)^k, \end{aligned} \tag{71}$$

then

$$e^{iy \hat{p}_x} f(x) e^{-iy \hat{p}_x} 1 = \sum_{k=0}^{\infty} c_k (x+y)^k = f(x+y), \tag{72}$$

as we wanted to prove.

Now consider the Bessel function J_n evaluated at $x+y$. From expression (69) we have

$$J_n(x+y) = e^{iy \hat{p}_x} J_n(x) e^{-iy \hat{p}_x} 1, \tag{73}$$

because $e^{-iy \hat{p}_x} 1 = 1$, and developing the first exponential in Taylor series, we obtain

$$J_n(x+y) = \sum_{m=0}^{\infty} \frac{y^m}{m!} \frac{d^m}{dx^m} J_n(x). \tag{74}$$

To calculate the m -derivative of J_n , we use the integral representation (66) to write

$$\frac{d^m}{dx^m} J_n(x) = i^m \frac{1}{2\pi} \int_{-\pi}^{\pi} \sin^m \tau e^{-i(n\tau - x \sin \tau)} d\tau, \tag{75}$$

substituting $\sin \tau = (e^{i\tau} - e^{-i\tau})/2i$, and using the binomial expansion,

$$\begin{aligned} \frac{d^m}{dx^m} J_n(x) &= \frac{1}{2^{m+1} \pi} \sum_{k=0}^m (-1)^k \binom{m}{k} \int_{-\pi}^{\pi} e^{i(m-k)\tau} e^{-ik\tau} e^{-i(n\tau - x \sin \tau)} d\tau \\ &= \frac{1}{2^m} \frac{1}{2\pi} \sum_{k=0}^m (-1)^k \binom{m}{k} \int_{-\pi}^{\pi} e^{-i[(n-m+2k)\tau - x \sin \tau]} d\tau, \end{aligned} \tag{76}$$

and therefore, using again the integral representation (66), we obtain

$$\frac{d^m}{dx^m} J_n(x) = \frac{1}{2^m} \sum_{k=0}^m (-1)^k \binom{m}{k} J_{n-m+2k}(x). \tag{77}$$

Substituting this last expression in equation (74), we obtain (we have taken the sum up to infinite as we add only zeros)

$$J_n(x+y) = \sum_{m=0}^{\infty} \frac{y^m}{m!} \frac{1}{2^m} \sum_{k=0}^{\infty} (-1)^k \binom{m}{k} J_{n-m+2k}(x). \quad (78)$$

We now change the order of summation and start the second sum at $m = k$ (because from $m < k$ all the terms are zero)

$$J_n(x+y) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \sum_{m=k}^{\infty} \frac{y^m}{2^m(m-k)!} J_{n-m+2k}(x). \quad (79)$$

We do now $j = m - 2k$ and obtain

$$J_n(x+y) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \sum_{j=-k}^{\infty} \frac{y^{j+2k}}{2^{j+2k}(j+k)!} J_{n-j}(x), \quad (80)$$

take the second sum from minus infinite, and exchange the order of the sums

$$\begin{aligned} J_n(x+y) &= \sum_{j=-\infty}^{\infty} J_{n-j}(x) \sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \frac{y^{j+2k}}{2^{j+2k}(m+k)!} \\ &= \sum_{j=-\infty}^{\infty} J_{n-j}(x) J_j(y). \end{aligned} \quad (81)$$

The final expression

$$J_n(x+y) = \sum_{k=-\infty}^{\infty} J_{n-k}(x) J_k(y) \quad (82)$$

is known as the *addition formula* for the Bessel functions of the first kind of integer order.

5 Conclusions

We have shown how to apply some of the formalism of operator theory to some special functions, namely, Hermite and Laguerre polynomials and Bessel functions. We have applied them to obtain some known series of functions, such as the addition formula for Bessel functions, and some, to our knowledge, not known series, such as the sum of even Hermite polynomials.

References

- [1] H.M. Moya-Cessa. New expressions for Laguerre and Hermite polynomials. arXiv:0809.2259v1 [math-ph], 2008.
- [2] G Dattoli, Generalized polynomials, operational identities and their applications, Journal of Computational and Applied mathematics, 118, 111-123, Elsevier (2000).
- [3] R.P. Soni W. Magnus, F. Oberhettinger. Formulas and Theorems for the Special Functions of Mathematical Physics. Springer-Verlag, 1966.

- [4] M. Abramowitz and I.A. Stegun. Handbook of Mathematical Functions With Formulas, Graphs and Mathematical Tables. National Bureau of Standards Applied Mathematics, 197
- [5] N. N. Lebedev. Special functions and their applications. Prentice-Hall, 1965.
- [6] Hans J. Weber & Frank E. Harris George B. Arfken. Mathematical Methods for Physicists. A Comprehensive Guide. Seventh edition. Elsevier, 2013.
- [7] I.S. Gradshteyn and I.M. Ryzhik. Table of Integrals, Series, and Products. Seventh edition. Elsevier, 2007.
- [8] William H. Louisell. Quantum Statistical Properties of Radiation. John Wiley & Sons, 1973.
- [9] Hector M. Moya-Cessa and Francisco Soto Eguibar. Differential Equations: An Operational Approach. Rinton Press, 2011.
- [10] Hector M. Moya-Cessa and Francisco Soto Eguibar. Introduction to Quantum Optics. Rinton Press, 2011.
- [11] Christopher Gerry and Peter Knight. Introductory Quantum Optics. Cambridge University Press, 2005.
- [12] I.X..Wong & D.R. Guo. Special functions. World Scientific, 1989.
- [13] H. Moya-Cessa and A. Vidiella-Barranco. On the interaction of two-level atoms with superpositions of coherent states of light. J. of Mod. Optics, 42:15471552, 1995.
- [14] Yury A. Brychkov. Handbook of Special Functions Derivatives, Integrals, Series and Other Formulas. CRC Press, 2008.
- [15] G. N. Watson. A treatise on the theory of Bessel functions. Cambridge University Press, 1995.
- [16] B.G. Korenev. Bessel functions and their Applications. Taylor & Francis, 2002.
- [17] F. Bowman. Introduction to Bessel functions. Dover Publications Inc, 1958.



Francisco Soto-Eguibar

is Titular Researcher at the Optics Department at the National Institute of Astrophysics, Optics and Electronics (INAOE) in Puebla, Mexico. He received his PhD degree in Physics at the National University of Mexico, in Mexico City in

1982. His research interest is in quantum optics and in the foundations of quantum mechanics.



Héctor M. Moya-Cessa obtained his PhD at Imperial College in 1993 and since then he is a researcher/lecturer at the National Institute of Astrophysics, Optics and Electronics (INAOE) in Puebla, Mexico where he works on Quantum Optics.

He is fellow of the Alexander von Humboldt Foundation and Regular Associate of the international Centre for Theoretical Physics (Trieste, Italy). INAOE in Puebla, Mexico.