

# Deriving Higher Order Rogue Waves Solutions in the Nonlinear System Using a Symbolic Computation Approach

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**Abstract:** In this paper, we consider a novel multiple rogue wave solutions of a (3 +1)-dimensional Kudryashov-Sinelshchikov equation by using a symbolic computation approach. Some higher order rogue wave solutions of a (3 +1)-dimensional Kudryashov-Sinelshchikov equation are presented. Some properties of the multi-rogue waves and their collision structures are given through numerical examples. Finally, some expected applications and extension have been mentioned.

**Keywords:** (3 +1)-dimensional Kudryashov-Sinelshchikov equation, Bilinear form, Rogue waves

## 1 Introduction

Rogue waves observed in the water tanks [1,2], deep ocean [3,4], and optical fibers [5,6]. To construct rogue wave solutions, many several methods have been proposed, including the inverse scattering method [7], the Hirota bilinear method [8], Darboux transformation method [9] and Backlund transformation method [10] and so on. Recently [11,12,13], the first order rogue wave and rational solutions to some (3+1) and (2+1)-dimensional systems are constructed by the symbolic computation approach.

In this paper, we mainly focus on a (3 + 1)-dimensional Kudryashov-Sinelshchikov equation. The description of the proposed method in [14] is given in section 2, the bilinear form and multi rogue waves for (3+1)-dimensional Kudryashov-Sinelshchikov equation derived in section 3, the first-order, second-order and third-order rogue wave solutions for (3+1)-dimensional Kudryashov-Sinelshchikov equation derived in section 4-6. Finally, conclusions are given in Section 7.

## 2 Analytical method

Suppose the nonlinear partial differential equations (NLEEs), take the form

$$H(u, u_t, u_x, u_y, u_z, u_{xt}, u_{xy}, u_{xz}, \dots) = 0, \quad (1)$$

where  $H$  is a polynomial in unknown function  $u(x, y, z, t)$  and its derivatives.

The main steps for the used method are consists of the following steps

**Step 1.** According the transformation to the Painlevé analysis

$$u(x, y, z, t) = u(f), \quad (2)$$

is made by a dependent variable function  $f$ .

**Step 2.** The NLEEs(1) using (2) converted into Hirota's bilinear form

$$F(D_\zeta, D_z) = 0, \quad (3)$$

where  $\zeta = x + y - et$ ,  $e$  is a real parameter and  $e$  is a wave speed D-operator [15] is defined by

$$D_x^k D_y^m D_z^n D_t^l f(x, y, z, t) \cdot g(x, y, z, t) = \left(\frac{\partial}{\partial x} - \frac{\partial}{\partial x'}\right)^k \left(\frac{\partial}{\partial y} - \frac{\partial}{\partial y'}\right)^m \left(\frac{\partial}{\partial z} - \frac{\partial}{\partial z'}\right)^n$$

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$$\left(\frac{\partial}{\partial t} - \frac{\partial}{\partial t'}\right)^l [f(x, y, z, t)g(x', y', z', t')] \big|_{x'=x, y'=y, z'=z, t'=t} \quad (4)$$

**Step 3.** Suppose

$$F = G_{n+1}(\zeta, z; \alpha, \beta) = F_{n+1}(\zeta, z) + 2\alpha z P_n(\zeta, z) + 2\beta(\zeta) Q_n(\zeta, z) + (\alpha^2 + \beta^2) F_{n-1}(\zeta, z), \quad (5)$$

$$\begin{aligned} F_n(\zeta, z; \alpha, \beta) &= \\ &\sum_{k=0}^{\frac{1}{2}n(n+1)} \left( \sum_{i=0}^k a_{n(n+1)-2k, 2i} z^{2i} \zeta^{n(n+1)-2k} \right), \\ P_n(\zeta, z) &= \\ &\sum_{k=0}^{\frac{1}{2}n(n+1)} \left( \sum_{i=0}^k b_{n(n+1)-2k, 2i} z^{2i} \zeta^{n(n+1)-2k} \right), \\ Q_n(\zeta, z) &= \\ &\sum_{k=0}^{\frac{1}{2}n(n+1)} \left( \sum_{i=0}^k c_{n(n+1)-2k, 2i} z^{2i} \zeta^{n(n+1)-2k} \right), \end{aligned} \quad (6)$$

where  $F_0 = 1, F_{-1} = P_0 = Q_0 = 0$ ,  $\alpha, \beta, a_{m,l}, b_{m,l}$  and  $c_{m,l}$ , ( $m, l = 0, 2, 4, \dots, n(n+1)$ ) are real numbers and  $\alpha, \beta$  are used to control the rogue-wave center.

**Step 4.** The systems obtained Solved by inserting (5) into (2) and equating all the coefficients of  $\zeta$  and  $z$  to zero we can obtain a system of polynomials and using the symbolic software Maple to solve the system.

**Step 5.** Substituting the values of  $a_{m,l}, b_{m,l}$  and  $c_{m,l}$  into (1) to get the multi rogue wave solutions in terms of  $\zeta$  and  $z$ .

### 3 The bilinear form for (3 + 1)-dimensional Kudryashov-Sinelshchikov equation

In this section, we consider the (3 + 1)-dimensional Kudryashov-Sinelshchikov equation [16]

$$(u_t + f_1 u u_x + f_2 u_{xxx})_x + f_3 u_{yy} + f_4 u_{zz} = 0, \quad (7)$$

where  $u(x, y, z, t)$  is a differentiable function and  $f_1, f_2, f_3, f_4$  are arbitrary constants.

To derive the rogue-waves solutions for (7), we find the Hirota bilinear form by setting  $\zeta = x + y - et$ , to obtain the ordinary differential equation (ODE) for (7) as follows

$$f_2 u_{\zeta\zeta\zeta\zeta} + (f_1 u + f_3 - e) u_{\zeta\zeta} + f_1 u_{\zeta}^2 + f_4 u_{zz} = 0. \quad (8)$$

Assume that the variable transformation has the following form

$$u = u_0 + \frac{12f_2}{f_1} (\ln F)_{\zeta\zeta}. \quad (9)$$

Then the Hirota bilinear form for (7) can be obtained by inserting (9) into (8) as

$$(f_2 D_{\zeta}^4 + (f_1 u_0 + f_3 - e) D_{\zeta}^2 + f_4 D_z^2) F \cdot F = 0. \quad (10)$$

The multi rogue wave solutions for the for a (3 + 1)-dimensional Kudryashov-Sinelshchikov equation (7) can be obtained as follows

### 4 First-order rogue waves $n = 0$

In this case, we select

$$F = G_1 = a_{2,0} \zeta^2 + a_{0,2} z^2 + a_{0,0}. \quad (11)$$

Let  $a_{2,0} = 1$  without loss of generality. Substituting (11) into (10) and setting the coefficients of all powers of  $\zeta$  and  $z$  to zero, we can obtain the coefficients  $a_{0,0}$  and  $a_{0,2}$  as follows

$$a_{0,0} = \frac{3f_2}{-u_0 f_1 + e - f_3}, a_{0,2} = -\frac{-u_0 f_1 + e - f_3}{f_4} \quad (12)$$

Inserting (12) in (11), then the first-order rogue waves for Eq. (7) takes the form

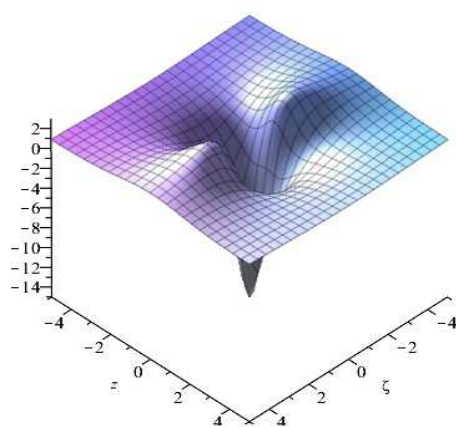
$$u = u_0 + \frac{12f_2}{f_1} (\ln F)_{\zeta\zeta}. \quad (13)$$

where

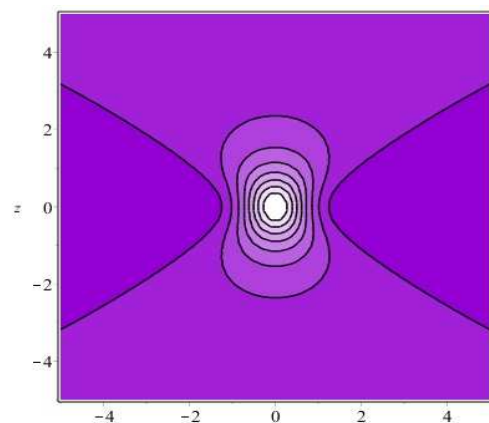
$$F = -\frac{(-u_0 f_1 + e - f_3)(z - \alpha)^2}{f_4} + (\zeta - \beta)^2 - \frac{3f_2}{u_0 f_1 - e + f_3}. \quad (14)$$

The first-order rogue wave solutions (13) when  $\alpha = \beta = 0$  are shown in Fig.1, has three center (0,0) and  $(\mp 3\sqrt{\frac{f_2}{-u_0 f_1 + e - f_3}}, 0)$  in three-dimensional, contour plot and the corresponding density plot, we notice that in this case, we have one peaks, the first-order rogue wave has the minimum amplitude  $u_0 + \frac{-8u_0 f_1 + 8e - 8f_3}{f_1}$  at (0,0) and maximal amplitude  $u_0 - \frac{-u_0 f_1 + e - f_3}{f_1}$  at  $(\mp 3\sqrt{\frac{f_2}{-u_0 f_1 + e - f_3}}, 0)$  when  $f_2 > 0, f_1 u_0 + f_3 < e$ . But the first-order rogue wave solutions (14) at  $\alpha = -2, \beta = -2$  the center of rogue wave will be  $(-2, -2)$  and  $(\frac{2U_0 f_1 + 3\sqrt{-u_0 f_1 f_2 + e f_2 - f_2 f_3 - 2e - 2f_3}}{-u_0 f_1 + e - f_3}, -2)$  as shown in Fig.2, moreover, the minimal and maximal amplitudes also change into  $u_0 + \frac{-8u_0 f_1 + 8e - 8f_3}{f_1}$  and  $-\frac{2u_0 f_1 + e - f_3}{f_1}$ , respectively.

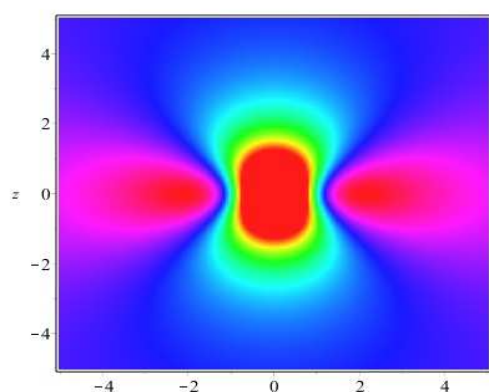
Fig.1 The first-order rogue wave solution (13) propagation in three-dimensional plot (a), contour plot (b) and density plot (c) with  $\alpha = \beta = 0$



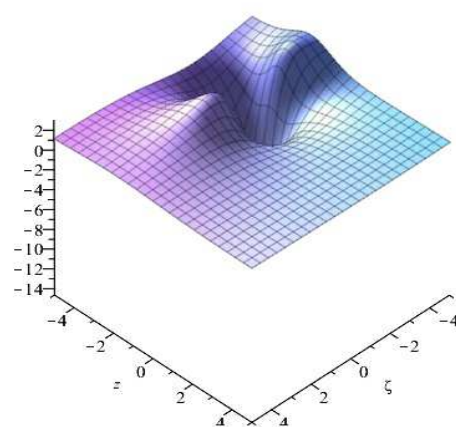
(a)



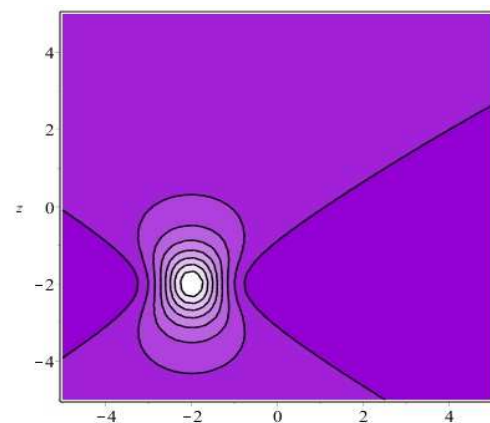
(b)



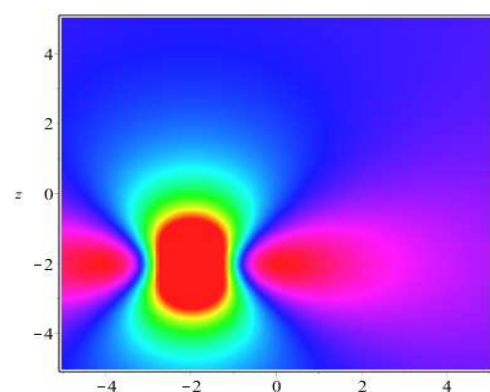
**Fig. 1:** The first-order rogue wave solution (13) propagation in three-dimensional plot (a), contour plot (b) and density plot (c) with  $\alpha = \beta = 0$



(a)



(b)



**Fig. 2:** The first-order rogue wave solution (13) propagation in three-dimensional plot (a), contour plot (b) and density plot (c) with  $\alpha = \beta = -2$

## 5 Second-order rogue waves $n=1$

In this section, we will find the second-order rogue wave of Eq. (7) by setting  $n = 1$  in Eq. (5) as follows

$$\begin{cases} F = G_2(\zeta, z; \alpha, \beta) = F_2(\zeta, z) + 2\alpha z P_1(\zeta, z) + \\ 2\beta \zeta Q_1(\zeta, z) + (\alpha^2 + \beta^2) F_0(\zeta, z) = \\ a_{6,0} \zeta^6 + (z^2 a_{4,2} + a_{4,0}) \zeta^4 + 2\zeta^3 \beta c_{2,0} + \\ (z^4 a_{2,4} + 2\alpha z b_{2,0} + z^2 a_{2,2} + a_{2,0}) \zeta^2 + 2\beta \\ (c_{0,2} z^2 + c_{0,0}) \zeta + a_{0,6} z^6 + a_{0,4} z^4 + 2\alpha z^3 b_{0,2} \\ + a_{0,2} z^2 + 2\alpha z b_{0,0} + a_{0,0} (\alpha^2 + \beta^2 + 1). \end{cases} \quad (15)$$

Substituting (14) into (10) and setting all coefficients of all powers for  $\zeta$  and  $z$  to zero, we can get a set of parameters  $a_{m,l}, b_{m,l}, c_{m,l}, (m, l = 0, 2, 4, 6)$  as follows

$$\begin{aligned} a_{0,0} &= \frac{1}{(\alpha^2 + \beta^2 + 1)(-u_0 f_1 + e - f_3)^3} \\ &\quad (-9\beta^2(-u_0 f_1 + e - f_3)^3 c_{2,0}^2 + \\ &\quad \alpha^2 u_0^2 b_{2,0}^2 f_1^2 f_4 - 2b_{2,0}^2 \alpha^2 u_0 f_4 (e - f_3) f_1 \\ &\quad + \alpha^2 e^2 b_{2,0}^2 f_4 - 2\alpha^2 e b_{2,0}^2 f_3 f_4 + \\ &\quad \alpha^2 b_{2,0}^2 f_3^2 f_4 - 16875 f_2^3), \\ a_{0,2} &= -475 \frac{f_2^2}{f_4 (-u_0 f_1 + e - f_3)}, \\ a_{0,4} &= 17 \frac{(-u_0 f_1 + e - f_3) f_2}{f_4^2}, \\ a_{0,6} &= -\frac{(-u_0 f_1 + e - f_3)^3}{f_4^3}, \\ a_{2,0} &= -125 \frac{f_2^2}{(-u_0 f_1 + e - f_3)^2}, \\ a_{2,4} &= 3 \frac{(-u_0 f_1 + e - f_3)^2}{f_4^2}, \\ a_{4,0} &= 25 \frac{f_2}{-u_0 f_1 + e - f_3}, \\ a_{4,2} &= -3 \frac{-u_0 f_1 + e - f_3}{f_4}, \\ b_{0,0} &= \frac{f_4 b_{2,0} f_2}{-3u_0 f_1 + 3e - 3f_3}, \\ b_{0,2} &= \frac{b_{2,0} (-u_0 f_1 + e - f_3)}{3f_4}, \\ c_{0,0} &= -\frac{c_{2,0} f_2}{-u_0 f_1 + e - f_3}, \\ c_{0,2} &= 3 \frac{(-u_0 f_1 + e - f_3) c_{2,0}}{f_4}, \\ a_{2,2} &= -90 \frac{f_2}{f_4} \end{aligned} \quad (16)$$

where  $b_{2,0}$  and  $c_{2,0}$  is an arbitrary constant. The second-order rogue wave for Eq. (7) takes the form

$$u = u_0 + \frac{12f_2}{f_1} (\ln G_2(\zeta, z; \alpha, \beta)) \zeta \zeta. \quad (17)$$

In Fig 3.,4 show the second-order rogue waves for (13) for different values of  $\alpha, \beta$ , the second-order peak breaks apart and for sufficiently big parameters, where the set of three first order rogue waves forms and these centers is formed a triangle, this is called a rogue wave triplet.

## 6 Third-order rogue waves $n = 2$

The third-order rogue wave of Eq. (7) is obtained by taking  $n = 2$  in Eq.(5) as follows

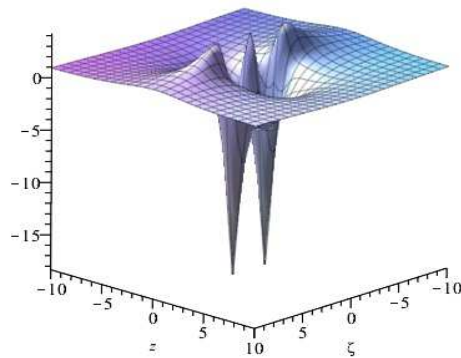
$$\begin{cases} F = G_3(\zeta, z; \alpha, \beta) \\ = F_3(\zeta, z) + 2\alpha z P_2(\zeta, z) + 2\beta \zeta Q_2(\zeta, z) \\ + (\alpha^2 + \beta^2) F_1(\zeta, z) \\ = a_{12,0} \zeta^{12} + a_{10,0} \zeta^{10} + a_{10,2} z^2 \zeta^{10} \\ + a_{8,0} \zeta^8 + a_{8,2} z^2 \zeta^8 + a_{8,4} z^4 \zeta^8 \\ + \zeta^6 + a_{6,2} z^2 \zeta^6 + a_{6,4} z^4 \zeta^6 \\ + a_{6,6} z^6 \zeta^6 + a_{4,0} \zeta^4 + a_{4,2} z^2 \zeta^4 \\ + a_{4,4} z^4 \zeta^4 + a_{4,6} z^6 \zeta^4 + a_{4,8} z^8 \zeta^4 \\ + a_{2,0} \zeta^2 + a_{2,2} z^2 \zeta^2 + a_{2,4} z^4 \zeta^2 + a_{2,6} z^6 \zeta^2 \\ + a_{2,8} z^8 \zeta^2 + a_{2,10} z^{10} \zeta^2 + 2\beta (c_{6,0} \zeta^6 + c_{4,2} z^2 \zeta^4 + c_{4,0} \zeta^4 \\ + c_{2,4} z^4 \zeta^2 + c_{2,2} z^2 \zeta^2 + c_{2,0} \zeta^2 \\ + c_{0,6} z^6 + c_{0,4} z^4 + c_{0,2} z^2 + c_{0,0}) (\zeta) + (\alpha^2 + \beta^2) \\ \times (a_{2,0} \zeta^2 + a_{0,2} z^2 + a_{0,0}) + a_{0,0} + 2\alpha z (b_{6,0} \zeta^6 \\ + b_{4,2} z^2 \zeta^4 + b_{4,0} \zeta^4 + b_{2,4} z^4 \zeta^2 \\ + b_{2,2} z^2 \zeta^2 + b_{2,0} \zeta^2 + b_{0,6} z^6 \\ + b_{0,4} z^4 + b_{0,2} z^2 + b_{0,0}) + a_{0,2} z^2 + a_{0,4} z^4 + a_{0,6} z^6 \\ + a_{0,8} z^8 + a_{0,10} z^{10} + a_{0,12} z^{12}. \end{cases} \quad (18)$$

Inserting (18) into (10), then setting all coefficients of all powers of  $\zeta$  and  $z$  to zero, we can get a set of parameters  $a_{m,l}, b_{m,l}, c_{m,l}, (m, l = 0, 2, 4, 6)$  where  $b_{2,0}$  and  $c_{4,0}$  are arbitrary constants. Then the third-order rogue wave solution for Eq. (7) takes the form

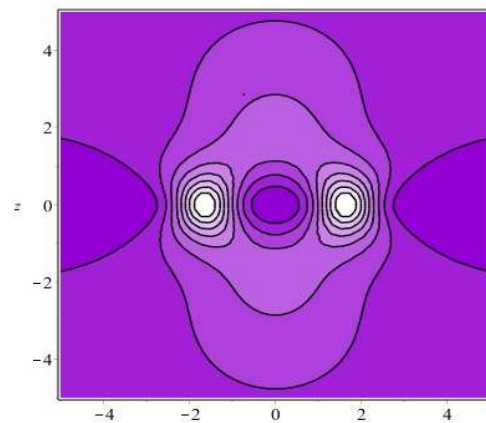
$$u = u_0 + \frac{12f_2}{f_1} (\ln G_3(\zeta, z; \alpha, \beta)) \zeta \zeta. \quad (19)$$

In Fig. 5,6. show the third-order rogue waves for (13) for big values of  $\alpha, \beta$ , the third-order rogue waves consists of five first-order rogue waves are located in the corners of a pentagon and other one sites in the center

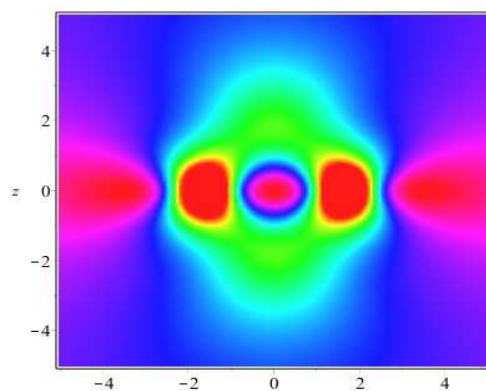
Fig.6. The third-order rogue wave solution (20) propagation in three-dimensional plot (a), contour plot (b) and density plot (c) with  $\alpha = \beta = 10^8$ . It is shown that multiple rogue waves solutions for a (3+1)-dimensional Kudryashov-Sinelshchikov equation can be used in different application such as matter-field interaction and Schrödinger equation [16]-[21].



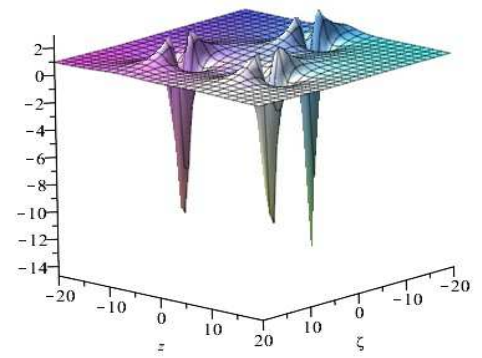
(a)



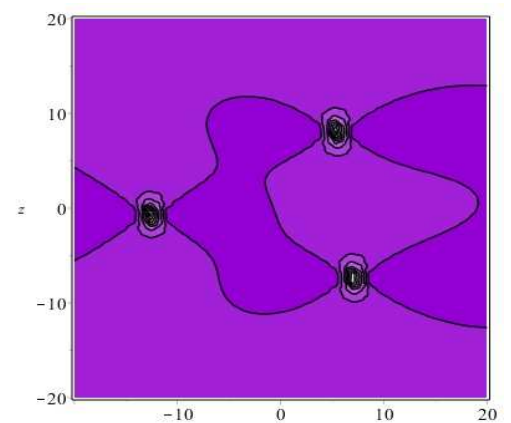
(b)



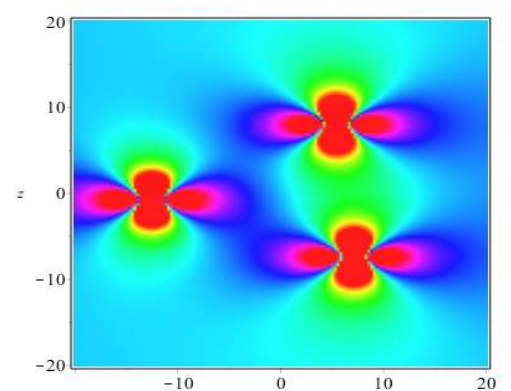
**Fig. 3:** The second-order rogue wave solution (17) propagation in three-dimensional plot (a), contour plot (b) and density plot (c) with  $\alpha = \beta = 0$



(a)

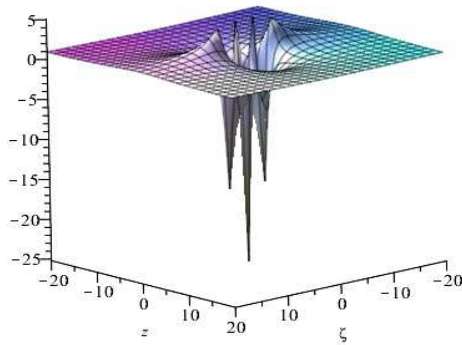


(b)

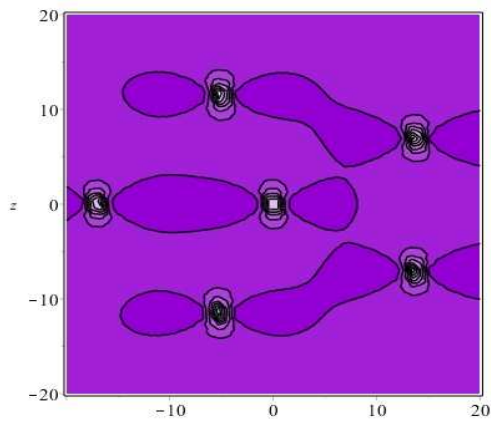


**Fig. 4:** Fig.4. The second-order rogue wave solution (17) propagation in three-dimensional plot (a), contour plot (b) and density plot (c) with  $\alpha = \beta = 2000$

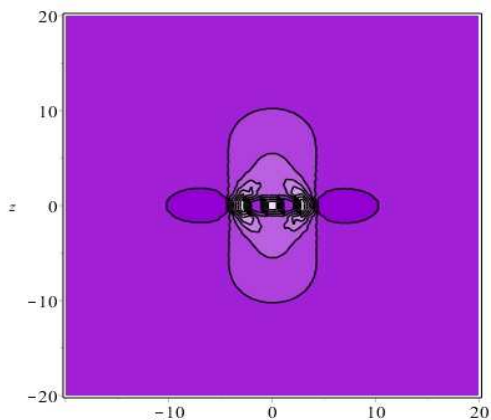




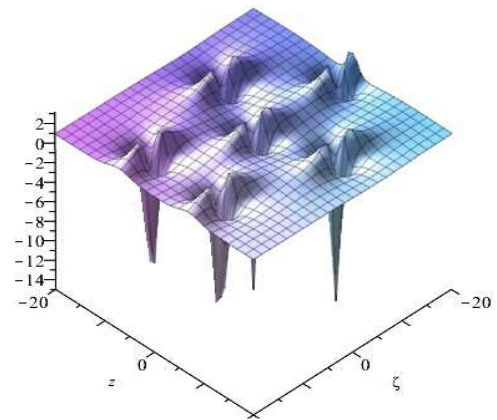
(a)



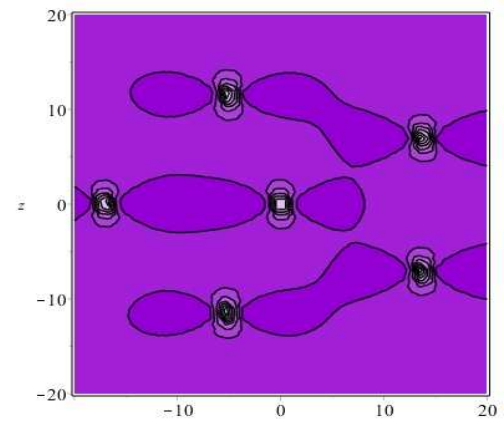
(b)



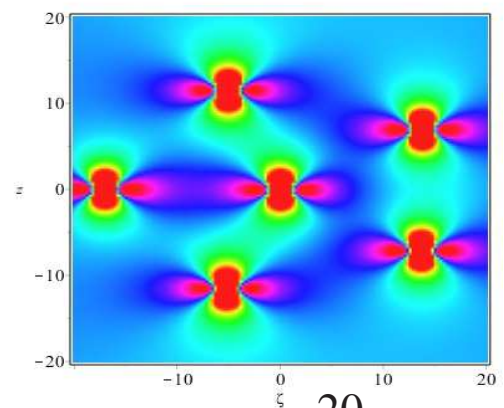
**Fig. 5:** The third-order rogue wave solution (20) propagation in three-dimensional plot (a), contour plot (b) and density plot (c) with  $\alpha = \beta = 0$



(a)



(b)



**Fig. 6:** The third-order rogue wave solution (20) propagation in three-dimensional plot (a), contour plot (b) and density plot (c) with  $\alpha = \beta = 10^8$

## 7 Conclusion

In this paper, we have discussed a novel multiple rogue wave solutions of a  $(3+1)$ -dimensional Kudryashov-Sinelshchikov equation contain two free parameters  $\alpha$  and  $\beta$ , which are used to control the center of the rogue waves. This exact solutions include higher order rogue wave solutions of a  $(3+1)$ -dimensional Kudryashov-Sinelshchikov equation.

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